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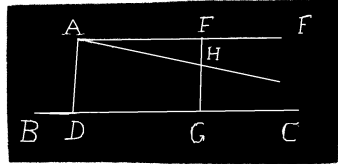
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to a given line. Draw AD perpendicular to BC , and through A draw AE perpendicular to AD . Then (1st and 2d) AE and DC are parallel; and the perpendicular FG equals AD , by (3d). Now suppose another line AH parallel to BC . Then HG equals AD or its equal FG . When HG equals FG , AH and AF coincide. Therefore, through a given point one line, and only one, can be drawn parallel to a given line.



The above demonstration may be made without the use of the word parallel. Thus: Through a given point one line, and only one, can be drawn equidistant from a given line.

With the figure drawn as in No. 3, begin the demonstration with the words: From *any* point F let fall the perpendicular &c., to prove the lines equidistant. Then with the same figure as in No. 4, and substituting the word equidistant for parallel, we have the demonstration.

CRADLE-ROCKING BY ELLEPTIC FUNCTIONS.

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After adopting the *gravitation-unit* of force, the equation of motion of the pendulum may be written $(h^2 + k^2)(W/g)(d^2\theta/dt^2) = -Wh \sin\theta \dots (1)$. Briefly making $(h + k^2/h) = l$ and $g/l = n^2$, we obtain from (1)

$$\frac{1}{2}(d\theta/dt)^2 = n^2(\text{vers } \alpha - \text{vers } \theta) \dots (2), \text{ in which, according to}$$

Sir William Thomson (Lord Kelvin), n is the *angular speed* of the pendulum, Divide *semicircularly* the pendulum-bob, turn downward the convex sides of these divisions centrally joined by a rectilinear axis of inappreciable length, and let the pendulum-rod bisect this rectilinear axis. In the position specified, these divisions constitute the *rockers* of an old-fashioned cradle; and this cradle we regard as placed upon a perfectly rough horizontal plane. Detaching the pendulum rod from the point of suspension, we have to consider *the rocking*, or the rolling oscillations on a horizontal plane, of a material body resting on a semicircular base. Let r = the radius of the equal semicircular rockers. Consider the *line* joining the points of tangency of the rockers with the horizontal plane, as the instantaneous axis of rotation; then, after obvious transformations, (2) becomes $\frac{1}{2}(r^2 - 2hr \cos\theta + h^2 + k^2)(d\theta/dt)^2 = gh(\text{vers } \alpha - \text{vers } \theta) \dots (3)$.

$$\begin{aligned}
 \therefore \left(\frac{dt}{d\theta}\right)^2 &= \frac{[(r-h)^2 + k^2] \cos^2 \frac{1}{2} \theta + [(r+h)^2 + k^2] \sin^2 \frac{1}{2} \theta}{4gh(\sin^2 \frac{1}{2} \alpha - \sin^2 \frac{1}{2} \theta)}, \\
 &= \frac{[(r-h)^2 + k^2] \cos^2 \frac{1}{2} \theta + [(r+h)^2 + k^2] \sin^2 \frac{1}{2} \theta}{4gh[\sin^2 \frac{1}{2} \alpha - (1 - \sin^2 \frac{1}{2} \alpha) \tan^2 \frac{1}{2} \theta]}, \\
 &= \frac{[(r-h)^2 + k^2] + [(r+h)^2 + k^2] \tan^2 \frac{1}{2} \theta}{4gh \cos^2 \frac{1}{2} \alpha (\tan^2 \frac{1}{2} \alpha - \tan^2 \frac{1}{2} \theta)} \dots (4).
 \end{aligned}$$

In order to transform (4), put $\tan \frac{1}{2} \theta = \tan \frac{1}{2} \alpha \cos \phi$; then differentiating,

$$\text{etc., } \left(\frac{d\theta}{d\phi}\right)^2 = 4 \left(\frac{\tan \frac{1}{2} \alpha \sin \phi}{1 + \tan^2 \frac{1}{2} \alpha \cos^2 \phi}\right)^2 = 4 \left(\frac{\sin \frac{1}{2} \alpha \cos \frac{1}{2} \alpha \sin \phi}{1 - \sin^2 \frac{1}{2} \alpha \sin^2 \phi}\right)^2 \dots (a).$$

$$\therefore \left(\frac{dt}{d\phi}\right)^2 = \frac{[(r-h)^2 + k^2] \cos^2 \frac{1}{2} \alpha + (1 - \sin^2 \phi) [(r+h)^2 + k^2] \sin^2 \frac{1}{2} \alpha}{gh(1 - \sin^2 \frac{1}{2} \alpha \sin^2 \phi)^2} \dots (5).$$

$$\text{Put } \kappa^2 = \frac{[(r+h)^2 + k^2] \sin^2 \frac{1}{2} \alpha}{[(r+h)^2 + k^2] \sin^2 \frac{1}{2} \alpha + [(r-h)^2 + k^2] \cos^2 \frac{1}{2} \alpha} \dots (b),$$

$$\text{and } \kappa'^2 = \frac{[(r-h)^2 + k^2] \cos^2 \frac{1}{2} \alpha}{[(r+h)^2 + k^2] \sin^2 \frac{1}{2} \alpha + [(r-h)^2 + k^2] \cos^2 \frac{1}{2} \alpha} \dots (c);$$

then will $\kappa^2 + \kappa'^2 = 1$. Make $n = \sqrt{g/r}$, and represent the denominators of (b) and (c) by M ; then (5) may be written

$$ndt = \sqrt{\left(\frac{M}{hr}\right)} \left[\frac{(1 - \kappa^2 \sin^2 \phi) d\phi}{(1 - \sin^2 \frac{1}{2} \alpha \sin^2 \phi) \sqrt{(1 - \kappa^2 \sin^2 \phi)}} \right] \dots (6).$$

According to the Jacobian system of notation as modified by Gudermann (*Theorie der Modular Functionen*), we have $\phi = \text{am } U$, and $\sin \frac{1}{2} \alpha = \kappa \text{sn } A$.

Since $\text{dn}^2 A + \kappa^2 \text{sn}^2 A = 1$, we obtain $\text{dn}^2 A = 1 - \kappa^2 \text{sn}^2 A = 1 - \kappa^2 (\sin^2 \frac{1}{2} \alpha / \kappa^2) = \cos^2 \frac{1}{2} \alpha$; also,

$$\text{sn } A = \sqrt{\left(\frac{[(r+h)^2 + k^2] \sin^2 \frac{1}{2} \alpha + [(r-h)^2 + k^2] \cos^2 \frac{1}{2} \alpha}{[(r+h)^2 + k^2]}\right)} \dots (d),$$

$$\text{and } \text{cn } A = \sqrt{(1 - \text{sn}^2 A)} = \frac{2 \cos \frac{1}{2} \alpha \sqrt{hr}}{\sqrt{[(r+h)^2 + k^2]}} \dots (e).$$

$$\therefore ndt = 2 \left(\frac{\text{sn } A \text{dn } A}{\text{cn } A}\right) \left[1 - \frac{\kappa^2 \text{cn}^2 A \text{sn}^2 U}{1 - \kappa^2 \text{sn}^2 A \text{sn}^2 U} dU \dots (7),\right.$$

and $nt = 2[(\text{sn } A \text{dn } A / \text{cn } A) U - \Pi(U, A)] \dots (8)$, while $\tan \frac{1}{2} \theta = \tan \frac{1}{2} \alpha \text{cn } U$.